# EVALUATING SADT BY ADVANCED KINETICS-BASED SIMULATION APPROACH

*B.* Roduit<sup>1\*</sup>, *P.* Folly<sup>2</sup>, *B.* Berger<sup>2</sup>, *J.* Mathieu<sup>2</sup>, *A.* Sarbach<sup>2</sup>, *H.* Andres<sup>3</sup>, *M.* Ramin<sup>3</sup> and *B.* Vogelsanger<sup>3</sup>

<sup>1</sup>AKTS AG Advanced Kinetics and Technology Solutions, TECHNOArk 3, 3960 Siders, Switzerland <sup>2</sup>armasuisse, Science and Technology, 3602 Thun, Switzerland <sup>3</sup>Nitrochemie Wimmis AG, 3752 Wimmis, Switzerland

Present study depicts the extension of the method of the application of the advanced kinetic description of the energetic materials decomposition by its combination with the exact heat balance carried out by numerical analysis and the determination of the self-accelerating decomposition temperature (SADT). Moreover, the additional parameters such as thermal conductivity of the self-reactive substances, the type of containers and insulation layers, and different temperature profiles of the surrounding environment were taken into consideration.

The results of DSC experiments carried out with different heating rates in the range of  $0.25-4^{\circ}$ C min<sup>-1</sup> were elaborated by the Thermokinetics software. The application the Thermal Safety software and the kinetics-based approach led to proper selection of experimental conditions for SADT testing. The applied approach enabled the simulation of such scenario as the thermal ignition of self-reactive chemicals conditioned previously for 12 h at 80°C and exposed later isothermally for 8 h to temperatures between 120–180°C. Described method can be used for analysis of possible development of runaway during storage or transport of dangerous goods (TDG) and containers, and subsequent choice of the conditions that can prevent an accident.

Keywords: cook-off, heat balance, isoconversional method, SADT, thermal decomposition kinetics, time to ignition, transport of dangerous goods, upscaling of DSC data

## Introduction

The self-accelerating decomposition temperature (SADT), introduced into the international practice by the United Nations, is an important parameter that characterizes thermal hazard under transport conditions of self-reactive substances. Important feature of the SADT is that it is not an intrinsic property of a substance but '...a measure of the combined effect of the ambient temperature, decomposition kinetics, packaging size and the heat transfer properties of the substance and its packaging'. The simulation of the SADT is a complicated task that requires the consideration of two issues: (i) the correct description of the kinetics of the process i.e. the temperature dependence of the reaction rate on the temperature and (ii) the precise link between kinetics and thermal behavior of the investigated materials. Both these tasks are strongly related because, without correct heat balance considerations, even the most exact evaluation only of the kinetics based on the experiments carried out in mg scale is of little value for the prediction of the real behaviour of materials in kg or Mg scale. On the other hand, the precise knowledge of the thermal properties of the

energetic materials when applied with kinetic parameters based on too simplified kinetic description of the decomposition (as e.g. still commonly used simple models such as 'zero-th' or 'first' order reaction) will not lead to the correct prediction of the SADT experiments as well.

Independent of the mass of the sample investigated in any thermoanalytical experiment the correct description of the behaviour of a reaction requires the knowledge of two important parameters (i) the kinetics of the investigated reaction and (ii) the heat balance of the system [1]. Depending on the mass of the sample both these parameters differently contribute to the reaction progress: (i) The kinetics of the process is always the same, because it is not depending on the mass of the substance, therefore the correct determination of the kinetic parameters allows simulation of the reaction progress under any temperature mode independent on the mass of investigated sample. (ii) On the other hand, the heat balance in the system strongly depends on the sample mass and therefore has to be considered in adiabatic and semiadiabatic conditions. One can distinguish following boundary cases:

<sup>\*</sup> Author for correspondence: b.roduit@akts.com

- in mg scale, i.e. under typical thermoanalytical experiments (DSC, HFC) one assumes the ideal exchange of the reaction heat with the environment, what, in turn, allows to carry on the experiments under isothermal conditions or with constant heating rates. The reaction heat does not influence the course of the reaction.
- in ton scale one assumes that the reactions proceeds in adiabatic conditions, without any heat exchange with surroundings, therefore all heat evolved during exothermic reactions remains in the sample increasing its temperature. Therefore the prediction of the reaction course has to take into account not only the kinetics of the reaction but also the temperature increase resulting from the self-heating.
- in kg scale the problem of the heat accumulation is more complicated because, depending on the thermal properties of the sample, its mass and type of the reactor (or container) a different amount of heat will be exchanged with the surroundings and some part of it will accumulate in the system. The heat balance in this situation is more complicated and generally can be correctly done only using numerical techniques such as finite element analysis, finite differences or finite volumes.

#### **Evaluation of the decomposition kinetics**

The evaluation of the kinetics of the decomposition of energetic materials is one of the main prerequisite necessary for the correct modelling of their properties. Generally, the kinetic parameters are calculated from the experimental data obtained by means of thermoanalyzers or calorimeters such as e.g. TG, DTA or DSC signals. For safety issues, the experiments should be run in 'closed high pressure sealed crucibles' (isochoric conditions) in order to consider the influence of the pressure on strongly pressure dependent reactions [2]. Independent of the technique applied, experimental the kinetic calculations require the dependence of the reaction extent  $\alpha$  on the time or temperature.

In DSC, the most commonly applied thermal analysis technique for examining energetic materials, the determination of the kinetic parameters from the recorded signal requires its integration in order to obtain the  $\alpha$ -time or temperature relationship necessary for kinetic calculations. Therefore, the DSC runs should be carried out till the end of the reaction and stopped slightly after this point, because the thermokinetic approach requires a baseline after the final peak for the determination of the heat of the reaction. The integration of DSC signal is influenced by the method chosen for the determination of the baseline.

Often applied the straight-line form of the baseline is incorrect [3]. The recorded signal depends not only on the heat of the reaction but is additionally affected by the change of the specific heat of the mixture reactant-products during the progress of the reaction.

With: B(t) – the baseline, S(t) – the differential signal (DSC), the reaction rate  $d\alpha/dt$  and progress  $\alpha(t)$  can be expressed as:

$$\frac{\mathrm{d}\alpha}{\mathrm{d}t} = \frac{S(t) - B(t)}{\int_{\mathrm{tod}}^{\mathrm{tend}} S(t) - B(t)]\mathrm{d}t}$$
(1)  
$$\alpha(t) = \frac{\int_{\mathrm{tod}}^{\mathrm{t}} (S(t) - B(t))\mathrm{d}t}{\int_{\mathrm{tod}}^{\mathrm{t}} S(t) - B(t)]\mathrm{d}t}$$
(2)

with

$$B(t) = [1 - \alpha(t)](a_1 + b_1 t) + \alpha(t)(a_2 + b_2 t) \quad (3)$$

where  $(a_1+b_1t)$  is a tangent at the beginning and  $(a_2+b_2t)$  a tangent at the end of the DSC signal S(t).

The construction of the baseline for the DSC signal obtained during decomposition of the single base propellant with a heating rate of 1 K min<sup>-1</sup> is depicted in the Fig. 1. It can be seen that the course of the baseline can significantly influence the determination of the heat of the reaction and the estimation of the  $\alpha$ -T dependence what, in turn, will significantly influence the determination of the kinetic parameters.

The very important feature of the AKTS-Thermokinetics Software [4] is the possibility of the optimization of the baseline for all experiments collected by different heating rates (or temperatures) so that the





random errors in the various baseline constructions for all heating rates will 'average themselves out'. This is because different baseline construction are just as likely to give a value of heat of reaction that is slightly too high as one that is too low.

When the decomposition follows a single kinetic model then the reaction can be described in terms of a single pair of Arrhenius parameters and the commonly used set of functions  $f(\alpha)$  reflecting the mechanism of the process. In such a case the dependence of the logarithm of the reaction rate over 1/T is linear with the constant slope m=E/R in full range of conversion degree  $\alpha$ . The reaction rate can be described by only one value of the activation energy E and one value of the pre-exponential factor A by the following expression:

$$\frac{\mathrm{d}\alpha}{\mathrm{d}t} = A \exp\left(-\frac{E}{RT(t)}\right) f(\alpha) \tag{4}$$

where *t* is time, T – temperature, R – the gas constant, E – the activation energy, A – the pre-exponential factor,  $\alpha$  is the fraction converted and  $f(\alpha)$  is a differential form of the conversion function depending on the reaction model.

However, the decomposition reactions are generally too complex to be described in terms of a single pair of Arrhenius parameters (A and E) and the commonly applied set of reaction models  $f(\alpha)$ . In general, decomposition reactions demonstrate profound multi-step characteristics. The assumption that the decomposition of an energetic material will obey a simple rate law is very rarely true. Moreover, the determination of the kinetic parameters from the single run recorded with one heating rate only (so called 'single curve' method) leads to erroneous results and according to the recent recommendations should not be applied anymore [5, 6].

As concluded in the International ICTAC Kinetics project [6–9], the proper calculation of the kinetics requires the series of non-isothermal measurements carried out at different heating rates. This procedure allows supplying the data set that generally contains the necessary amount of information required for full identification of the complexity of a process.

In the present paper the kinetic parameters have been calculated by the isoconversional method of Friedman [10] based on the calculation of E and Avalues at different degrees of conversion  $\alpha$  without assuming the form of  $f(\alpha)$  function, i.e. applying logarithmic form of the following reaction rate expression:

$$\frac{\mathrm{d}\alpha}{\mathrm{d}t_{\alpha}} = \left[A_{\alpha}f(\alpha)\right]\exp\left(-\frac{E_{\alpha}}{RT(t_{\alpha})}\right) \tag{5}$$

according to Friedman we obtain

$$\ln\left(\frac{\mathrm{d}\alpha}{\mathrm{d}t_{\alpha}}\right) = \ln[A_{\alpha}f(\alpha)] - \frac{E}{R}\frac{1}{T_{\alpha}} \tag{6}$$

where  $t_{\alpha}$ ,  $T_{\alpha}$ ,  $E_{\alpha}$  and  $A_{\alpha}$  are the time, temperature, apparent activation energy and preexponential factor, at conversion  $\alpha$ , respectively, and  $-E_{\alpha}/R$  and  $\ln[A_{\alpha}f(\alpha)]$  are the slope and the intercept with the vertical axis of the plot of  $\ln(d\alpha/dt_{\alpha})$  vs.  $1/T_{\alpha}$ .

It is then possible to make kinetic predictions at any temperature profile T(t), from the values of  $E_{\alpha}$ and  $[A_{\alpha}f(\alpha)]$  extracted directly from the Friedman method by the separation of the terms followed by an integration:

$$t_{\alpha} = \int_{0}^{t_{\alpha}} dt = \int_{0}^{\alpha} \frac{d\alpha}{[A_{\alpha}f(\alpha)]e^{-\frac{E_{\alpha}}{RT_{\alpha}}}}$$
(7)

No specification of the reaction model term,  $f(\alpha)$ , is necessary for the kinetic prediction since it requires (Eq. (7)), along with  $E_{\alpha}$ , only the product term,  $[A_{\alpha}f(\alpha)]$ , which is experimentally extracted from the kinetic experiment according to Eq. (6). However, nothing can be inferred about the pre-exponential factor,  $A_{\alpha}$ , unless  $f(\alpha)$  is assumed to have some particular form (first-order reaction, nucleation-growth, etc.). And when  $f(\alpha)$  is associated with a specific reaction model, still the experimentally extracted product term  $[A_{\alpha}f(\alpha)]$  remains unchanged. The AKTS-Thermokinetics Software and AKTS-Thermal Safety Software [4] take advantage of this unique property of the Friedman method of the isoconversional kinetics evaluation and enable determination and displaying, in arbitrarily chosen form, the apparent 'first-order pre-exponential factor',  $A_{\alpha}$  (Fig. 2).

The results of the determination of the kinetic parameters for a 155 mm artillery charge of Swiss army are depicted in Fig. 2. Figures 2a–c present the dependence of E and A on the reaction progress. Figures 2b–d show the comparison of the experimental and calculated courses of the reactions for different heating rates by applying the determined E and A values. Note the very good fit of calculated and experimental relationships.

The application of the advanced kinetics for the investigation of the properties of energetic materials was also recently presented by Burnham and Dinh [11].

#### Heat balance

To perform the exact heat balance the numerical techniques like finite element analysis or finite



Fig. 2 Kinetic analysis of 155 mm artillery charge of Swiss army: a, b - combustible cartridge, c, d - single-base propellant. a, c - activation energy E and pre-exponential factor A as a function of the reaction progress. b, d - normalized DSC-signals as a function of the temperature and heating rate (marked in °C min<sup>-1</sup> on the curves). Experimental data are depicted as symbols, solid lines represent the signals calculated on the basis of kinetic parameters determined by the Friedman method

conduction to its boundaries. Mathematically, using Fourier's law of heat conduction, we can derive from the heat equation:  $\frac{\partial T}{\partial t} = \frac{\lambda}{\rho C_{\rm P}} \nabla^2 T + \frac{1}{\rho C_{\rm P}} q_{\rm r}$ 

time and space.

where  $\lambda$ ,  $\rho$ ,  $C_{\rm P}$ , T,  $q_{\rm r}$  mean: thermal conductivity, density, specific heat, temperature and the power generated per unit volume by the decomposition reaction, respectively. With

differences or volumes can be applied. The sample is

virtually divided into the set of adjoining elements

(Fig. 3). These elements are organized in a virtual

mesh and described by the advanced thermokinetics

based on the Friedman analysis in each node of the

depends on the rate of heat generation, its capacity to

store the part of this heat, and the rate of heat

When the heat is transferred to the surrounding environment the temperature profile within a body

$$q_{\rm r} = \rho \Delta H_{\rm r} \frac{\mathrm{d}\alpha}{\mathrm{d}t} \tag{9}$$

(8)

after considering cylindrical coordinates and some simplifications,

$$\frac{\partial T}{\partial r} \rangle \rangle \frac{\partial T}{\partial z} \Rightarrow \frac{\partial^2 T}{\partial r^2} \rangle \rangle \frac{\partial^2 T}{\partial z^2}$$
(10)

one can write

$$\frac{\partial T}{\partial t} = \frac{\lambda}{\rho C_{\rm P}} \left( \frac{\partial^2 T}{\partial r^2} + \frac{J}{r} \frac{\partial T}{\partial r} \right) + \Delta T_{\rm ad} \frac{\mathrm{d}\alpha}{\mathrm{d}t} \qquad (11)$$

where J is a geometry factor dependent on the type of the container: J=0 for the infinite plate, J=1 for the infinite cylinder and J=2 for the sphere,  $d\alpha/dt$  is the rate of the decomposition reaction expressed by the Arrhenius type equation as those applied in Friedman analysis (Eq. (5)) and  $\Delta T_{ad}$  is the adiabatic temperature rise expressed by the heat of reaction  $\Delta H_{\rm r}$  and the specific heat  $C_{\rm P}$ :  $\Delta T_{\rm ad} = -\Delta H_{\rm r}/C_{\rm P}$ .

## Simulation of properties of 155 mm artillery charge of Swiss army

The extension of the simulations under isothermal conditions on the prediction of the time to selfignition is illustrated by the results depicting the investigation of a 155 mm artillery charge used in Swiss army. During modeling the geometry, dimension of the container and, additionally, the amount, properties and thickness of the layers of different materials used for the construction of the ammunition container have been taken into account. Application of the appropriate decomposition kinetics (results

#### EVALUATING SADT BY ADVANCED KINETICS-BASED SIMULATION APPROACH



Fig. 3 Generalized heat balance over a container and a volume element; a – kinetic parameters calculated from the DSC measurements, independent of the sample mass, enable the determination of the reaction rate required for the heat balance; b – heat balance depends on the sample mass and has to be calculated by the numerical techniques

Table 1 The properties of the components of 155 mm artillery charge

	Container wall	Protection material	Combustible cartridge case	Single-base propellant
Layer thickness/mm	0.7	9	3	64.5
Thermal conductivity $\lambda/W \text{ cm}^{-1} \text{ K}^{-1}$	0.163	0.0023	0.002	0.001
Density $\rho/g \text{ cm}^{-3}$	8	0.4	1.38	1.38
Specific heat $C_{\rm P}/{\rm J~g}^{-1}~{\rm K}^{-1}$	0.5	2.3	1.5	1.5
Heat of reaction $\Delta H_r/J \text{ g}^{-1}$	_	_	$-2768\pm256$	$-3579 \pm 350.6$

presented in Fig. 2) and proper heat balance calculated by finite element analysis enabled the determination of the effect of scale and geometry of the container as well as the heat transfer, thermal conductivity and surrounding temperature on the heat accumulation in the sample. The simulation results together with the experimental data are depicted in Fig. 4. The simulation has been done for the sample in the form of cylinder containing four layers of the following materials possessing significantly different thermal properties: single-base propellant, combustible cartridge case, protection material and steel container. The properties of these materials, required for the modeling, are summarized in Table 1.

Additionally to the isothermal simulations and investigations of the properties of the energetic materials also the commonly applied 'slow cook-off' experiment (heating rate  $3.3^{\circ}$ C h<sup>-1</sup>, initial temperature

of the charge 40°C, 6 h) has been carried out with the 4.5 kg artillery charge. The experimental data and results of the simulation of the slow cook-off investigation are presented in Fig. 5. As for the isothermal simulation, the heat balance calculated by finite element analysis was applied together with the advanced kinetic description of the reaction. The experimentally determined ignition temperature of artillery charge amounted to 137°C (Fig. 5a). The predicted ignition temperature of 138°C (Fig. 5b) was in a very good agreement with the experimental value.

## Self-accelerating decomposition temperature (SADT)

SADT is an important parameter that characterizes thermal hazard under transport conditions of self-



**Fig. 4** Single-base propellant for a 155 mm artillery charge applied in Swiss army: the dependence of the calculated time to ignition on the temperature under non-adiabatic conditions



**Fig. 5** Slow cook-off of the 155 mm artillery charge a – experiment and b – simulation. The predicted temperature of explosion 137°C was in good agreement with the slow cook-off experimental value of 138°C

reactive substances. The SADT has been introduced into the international practice by the regulations of the United Nations presented in 'Recommendations on the Transport of Dangerous Goods, Manual of Tests and Criteria' (TDG) [12]. The Globally Harmonized System (GHS) [13] has inherited the SADT as a classification criterion for self-reactive substances. According to the Recommendations on TDG the SADT is defined as 'the lowest temperature at which self-accelerating decomposition may occur with a substance in the packaging as used in transport'. Important feature of the SADT is that it is not an intrinsic property of a substance but '...a measure of the combined effect of the ambient temperature, decomposition kinetics, packaging size and the heat transfer properties of the substance and its packaging' [12].

The Manual of Tests and Criteria of the United Nations on the transport of dangerous goods and on the globally harmonized system of classification and labelling of chemicals indicates that the characterization of the materials is based on the heat accumulation storage tests. One can find the following definitions: (*i*) SADT is the lowest environment temperature at which overheat in the middle of the specific commercial packaging exceeds  $6^{\circ}$ C ( $\Delta T_6$ ) after a lapse of the period of seven days (168 h) or less. This period is measured from the time when the packaging center temperature reaches 2°C below the surrounding temperature. (*ii*) SADT is the critical ambient temperature rounded to the next higher multiple of 5°C.

The first definition is based on two essential parameters – maximal permissible overheating temperature and minimal acceptable induction period. The second definition considers only one parameter: the critical ambient temperature of thermal runaway rounded to the next higher multiple of 5°C without any fixed transportation time in the definition.

## Up-scaling of DSC data for the determination of the SADT

Having the advanced kinetic description performed by AKTS-Thermokinetics Software, the numerical techniques implemented in AKTS-Thermal Safety Software enable to simulate the heat conduction properties of the energetic material stored in a specific container. Thus, applying previously determined kinetic parameters and heat balance, it was possible to determine precisely the SADT for the 155 mm artillery charge as illustrated in Fig. 6. For the determination of the SADT we had to specify the properties of the containers. The simulations have been done for a container containing four layers of the materials possessing significantly different thermal properties: single-base propellant, combustible cartridge case, protection material and steel container. The properties of these materials, required for the modelling, are summarized in Table 1. For the determination of the SADT in the current study a heat transfer coefficient  $U=30 \text{ Wm}^{-2} \text{ K}^{-1}$  has been chosen according to the EU norm recommendations EN1991-1-2 (2002). For each simulation, the initial temperature in each layer is 20°C. The SADT based on the first definition (i) amounts to 100°C (Fig. 6a).



**Fig. 6** Determination of SADT and critical temperature of thermal runaway ( $T_{CR}$ ). a – based on the first definition (*i*) we obtain a SADT of 100°C; b – for the  $T_{CR}$  of 65°C the time to thermal runaway amounts to ca. 80 years; c – development of the temperatures of the container (bottom curve – the temperature of the wall, top curve – the temperature in the center) at the SADT (*ii*) of 70°C

The temperature course of the reaction reveals that it will probably proceed in a thermal runaway scenario soon after 8 days. Anyhow, based on the definition,  $\Delta T_6$  is reached after a lapse of about 6.7 days. This period is measured from the time when the packaging centre temperature reaches 2°C below the surrounding temperature.

Bearing in mind the above-mentioned SADT definitions, it is important to understand in more details how the 'SADT' correlates with the 'critical temperature of thermal runaway' ( $T_{CR}$ ). For a packaging of given size, the  $T_{CR}$  delimits the runaway and non-runaway domains of reaction proceeding and represents fundamental attribute of a thermal runaway. Due to the fact that simulation enables calculation of a possible thermal runaway as a function of any surrounding temperature one can deliberately avoid considering any fixed transportation time in the definition. It may be or 7 days as in the TDG definition (*i*) or any undefined time if the second definition (*ii*) is considered. Thus, Fig. 6b presents the

dependence of the thermal runaway time as a function of the ambient temperature calculated according to this second definition of the SADT. The critical ambient temperature  $T_{CR}$  for the examined cylinder (Fig. 6b) is 65°C, what, according to the second SADT definition (ii), results in the SADT temperature of 70°C. However, under these conditions the induction period is about 78 years! Particularly, if a selfaccelerating reaction proceeds in a substance, the critical temperature based on definition (ii) is reached after a period longer than 7 days and it is lower than the SADT based on definition (i). This is very often the case with autocatalytic reactions, for such selfaccelerating reactions, the SADT is always higher than the critical temperature and this difference may reach even 35°C as illustrated in the present study.

Note the similarity of Figs 4 and 6b presenting the dependence of the time to ignition on the temperature under non-adiabatic conditions. Figure 6c presents the change of the temperature in the container at the SADT (ii) of 70°C, the runaway process occurs after ca. 33 years. The evidence of the strong autocatalytic behaviour of the investigated energetic materials can be easily illustrated by displaying the reaction rate under isothermal conditions (Fig. 7). The decomposition reaction is strongly autocatalytic because the initial rate of the reaction is low with a long induction time under isothermal conditions. This observation is in accordance with other studies of isothermal experiments at 100 and 110°C [14] showing experimentally that propellants do tend to have a rapid decomposition after a long induction time at modest temperatures during isothermal heating. Note that if we would be facing a non-self-accelerating reaction (such as 'n<sup>th</sup>-order' decomposition reaction types), the SADT would be equal or slightly lower than critical temperature and it would be reached after the period shorter than 7 days.

Thus the kinetic based method can be used for analysis of possible development of runaway during storage or transport of dangerous goods (TDG) and containers as specified in the first definition (i) of the SADT considering a period of transport of seven days or less. However, the method can also be applied for finding out the exact critical ambient temperature leading to a thermal runaway as specified in the second definition (ii) without specifying any fixed transportation time. Therefore, by examining the type of the decomposition kinetics, the kinetic based approach enables to give always the most critical value for the SADT. In addition to the experiments, it can be therefore considered as a reliable and quick method to increase safety during transport and storage of dangerous chemicals.



**Fig. 7** Simulated reaction rates of the single-base propellant under isothermal conditions (100, 110 and 120°C). The reaction is strongly autocatalytic because the initiation rate of the reaction is low leading to a long induction time under isothermal conditions

Note that the presented approach can be easily extended to the simulation of the impact of the fire. In the current study a heat transfer coefficient U=30 W m<sup>-2</sup> K<sup>-1</sup> has been chosen according to the EU norm recommendations EN1991-1-2 (2002). In case of fire, this heat transfer coefficient can be adjusted correspondingly according to same EU norm. During fire, the heat transfer occurs by both convection and radiation. Generally, the radiation is more dominant than the convection after the very early stages of the fire. The thermal actions can be represented by the net heat flux to the surface of the container. On the fire exposed surfaces, the net heat flux  $Q_{\text{net}}$  considering heat transfer by convection and radiation can be determined by:

$$Q_{\rm net} = Q_{\rm net,c} + Q_{\rm net,r} \tag{12}$$

where  $Q_{\text{net,c}}$  is the net convective heat flux component as given in Eq. (13):

$$Q_{\rm net,c} = U_{\rm c}(T_{\rm g} - T_{\rm s})({\rm W m}^{-2})$$
 (13)

where  $T_{\rm g}$  – the surrounding gas temperature around the container (°C),  $T_{\rm s}$  – the surface temperature of the container (°C),  $Q_{\rm net,r}$  is the net radiative heat flux component that can be calculated using the Stefan–Boltzmann law:

$$Q_{\text{net,r}} = \Phi \varepsilon_{\text{s}} \varepsilon_{\text{r}} \sigma [(T_{\text{r}} + 273.15)^4 - (T_{\text{s}} + 273.15)^4] \quad (14)$$

where  $\varepsilon_r$  – the emissivity of the fire (=1.0),  $\varepsilon_s$  – the surface emissivity of the container,  $\Phi$  – the configuration factor (≤1.0),  $T_r$  – the effective radiation temperature of the fire environment (°C),  $T_s$  – the surface temperature of the container (°C),  $\sigma$  – the Stephan Boltzmann constant (=5.67·10<sup>-8</sup> W m<sup>-2</sup> K<sup>-4</sup>).

Using  $Q_{\text{net}}=Q_{\text{net,c}}+Q_{\text{net,r}}$  one can now generalize the kinetic based approach by estimating an equivalent heat transfer coefficient (convection and radiation) for any given surrounding temperature with

$$U_{\text{equivalent}} = U_{\text{c}} + U_{\text{r}} = U_{\text{c}} + \frac{Q_{\text{r,net}}}{(T_{\text{g}} - T_{\text{s}})}$$
(15)

## Conclusions

Independent of the mass of the sample investigated in any thermoanalytical experiment, the correct description of the time to thermal runaway of a decomposition reaction requires the knowledge of two important parameters (i) the kinetics of the investigated reaction and (ii) the heat balance of the system. Depending on the mass of the sample both these parameters differently contribute to the reaction progress. Therefore, the advanced method of the simulation of the thermal properties of the energetic materials consists in the determination of the kinetic parameters of the decomposition process which are applied later, together with the exact heat balance calculated by finite element analysis method, for the prediction of the thermal behaviour of the investigated material under any temperature mode. The results of the simulation of the time-to-ignition under non-adiabatic conditions and 'slow cook-off' experiment for a 155 mm artillery charge of Swiss army indicated the very good fit of the simulated and the experimental data. For the slow cook-off experiment the predicted temperature of the explosion was 137°C whereas the experimentally found temperature amounted to 138°C. During the simulation, many parameters influencing the heat transport such as dimensions and shape of the container, amount and thickness of the layers of the materials used for the construction of the ammunition, their specific heat and thermal conductivity have been taken into account. The possibility of considering all these factors enabled the determination of the SADT or, according to its both, recommended definitions or, under specific, arbitrarily chosen scenario. Additionally, the application of the AKTS Thermal Safety Software can be extended for the prediction of the time to thermal runaway in the thermal environment created by the fire at conditions included in EN Standard concerning fire resistance.

#### References

- B. Roduit, C. Borgeat, B. Berger, P. Folly, H. Andres, U. Schädeli and B. Vogelsanger, J. Therm. Anal. Cal., 85 (2006) 195.
- 2 Swiss Institute of Safety and Security: http://www.swissi.ch/index.cfm?rub=1010.
- 3 W. F. Hemminger and S. M. Sarge, J. Thermal Anal., 37 (1991) 1455.

- 4 Advanced Kinetics and Technology Solutions: http://www.akts.com (AKTS-Thermokinetics software and AKTS-Thermal Safety software).
- 5 M. E. Brown, J. Therm. Anal. Cal., 82 (2005) 665.
- 6 M. Maciejewski, Thermochim. Acta, 355 (2000) 145.
- 7 M. E. Brown, M. Maciejewski, S. Vyazovkin, R. Nomen, J. Sempere, A. Burnham, J. Opfermann, R. Strey, H. L. Anderson, A. Kemmler, R. Keuleers, J. Janssens, H. O. Desseyn, C. R. Li, T. B. Tang, B. Roduit, J. Malek and T. Mitsuhashi, Thermochim. Acta, 355 (2000) 125.
- 8 A. Burnham, Thermochim. Acta, 355 (2000) 165.
- 9 B. Roduit, Thermochim. Acta, 355 (2000) 171.
- 10 H. L. Friedman, J. Polym. Sci, Part C, Polymer Symposium (6PC), 183 (1964).

- 11 A. K. Burnham and L. N. Dinh, J. Therm. Anal. Cal., 89 (2007) 479.
- 12 2003, Recommendations on the Transport of Dangerous Goods, Manual of Tests and Criteria, 4 revised Ed., United Nations, ST/SG/AC.10/11/Rev.4 (United Nations, New York and Geneva).
- 13 2003, Globally Harmonized System of Classification and Labelling of Chemicals (GHS), United Nations, New York and Geneva.
- 14 B. Roduit, C. Borgeat, B. Berger, P. Folly, B. Alonso and J. N. Aebischer, J. Therm. Anal. Cal., 80 (2005) 91.

DOI: 10.1007/s10973-007-8865-2